

# Physics 625

## Fizeau

FIZEAU: Based on interference

Suppose there are two light sources emitting monochromatic light of the same frequency ( $\nu$ ). The sources are separated by a distance  $\gg \lambda$  ( $\lambda$  = wavelength of light), and the detector is far enough from the sources to treat the sources as plane waves. If  $\vec{E}_1$  and  $\vec{E}_2$  are the electric fields of the two light sources, the irradiance at detector point (P) is:

$$I = \langle \vec{E}_1^2 \rangle_T + \langle \vec{E}_2^2 \rangle_T + 2 \langle \vec{E}_1 \cdot \vec{E}_2 \rangle_T$$

What is  $\langle \rangle_T$ ?  $\langle f(t) \rangle_T = \frac{1}{T} \int_t^{t+T} f(t') dt'$

For us, the time interval (T) is  $\gg$  period of light ( $= \frac{1}{\nu}$ )

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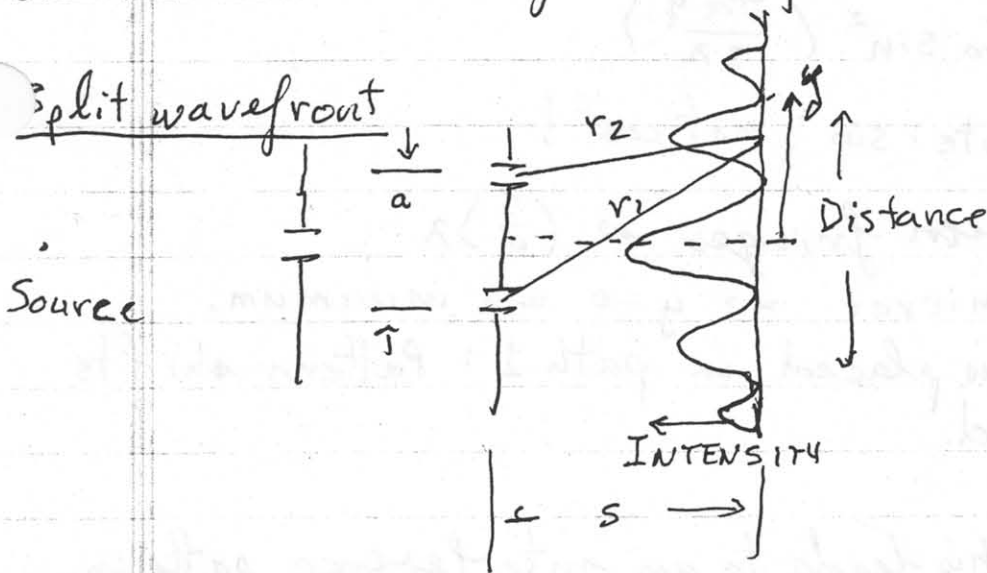
Naturally, if  $\vec{E}_1 \perp \vec{E}_2$ , there is no interference. When  $\vec{E}_1 \parallel \vec{E}_2$ , the interference term:

$$2 \langle \vec{E}_1 \cdot \vec{E}_2 \rangle_T = E_{01} E_{02} \cos(\delta)$$

$\delta = \vec{k}_1 \cdot \vec{r} - \vec{k}_2 \cdot \vec{r} + \epsilon_1 - \epsilon_2$  is the phase difference due to both the path length  $(\vec{k}_1 \cdot \vec{r} - \vec{k}_2 \cdot \vec{r})$  and initial phase difference  $(\epsilon_1 - \epsilon_2)$ .

continuous

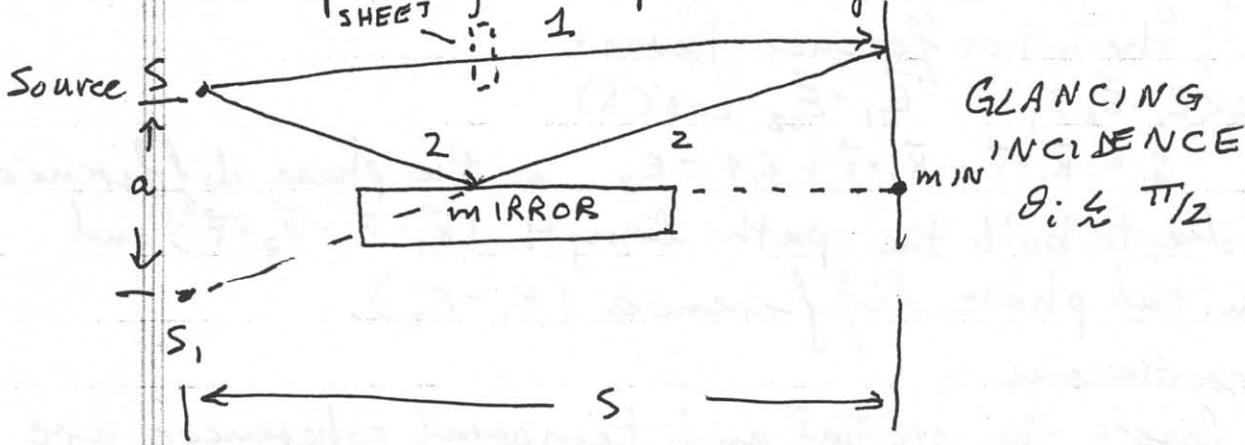
Now, with lasers the spatial and temporal coherences are long compared to the response time ( $\tau$ ) of the detectors. What about before 1948 - no lasers? How to get coherence needed for interference?



$$r_1 - r_2 \approx \frac{ya}{s}$$

$$\begin{aligned} \text{Interference intensity } I &= 4 I_0 \cos^2 \left( \frac{k(r_1 - r_2)}{2} \right) \\ &= 4 I_0 \cos^2 \left( \frac{ya\pi}{s\lambda} \right) \end{aligned}$$

Other wave-splitting example: Lloyd's mirror



At glancing incidence: Reflected beam has  $180^\circ$  phase shift

$$\Rightarrow \delta = k(r_1 - r_2) \pm \pi$$

$$\Rightarrow I = 4 I_0 \sin^2 \left( \frac{\pi a y}{s \lambda} \right)$$

Note:  $\sin^2$ , not  $\cos^2$ !

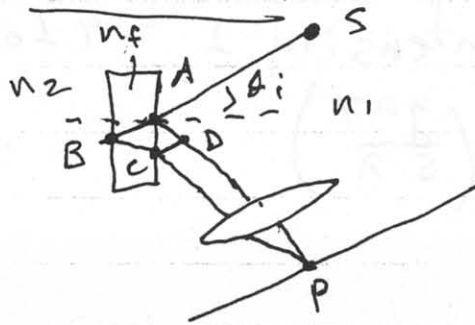
Spacing between fringes is  $(\frac{s}{a})\lambda$

Parallel to mirror  $\Rightarrow y=0 \Rightarrow$  minimum.

If a sheet is placed in path 1: Pattern shifts upward.

Split Amplitude: This leads to an interference pattern, to bright and dark fringes.

Equal inclination



$d =$  film thickness

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Difference in optical path length =  $\Lambda = n_f (\overline{AB}) + n_f (\overline{BC}) - n_i (\overline{AD})$

$$\overline{AB} = \overline{BC} = \frac{d}{\cos \theta_t}$$

$$\overline{AD} = \overline{AC} \sin \theta_i = \overline{AC} \left( \frac{n_f}{n_i} \right) \sin \theta_t$$

$$\overline{AC} = 2d \tan \theta_t$$

$$\Rightarrow \Lambda = 2n_f d \cos \theta_t$$

There is also a phase shift due to the reflections:

$0^\circ \leq \theta_i \leq 30^\circ$  two beams have a  $(\pi)$  relative phase shift

So if the phase shift due to the path difference is  $(k_0 \Lambda)$ , the total phase shift is:

$$\delta = k_0 \Lambda \pm \pi = \frac{4\pi d}{\lambda_0} (n_f^2 - n_i^2 \sin^2 \theta_i)^{1/2} \pm \pi$$

Interference maximum: when  $\delta = 2\pi m$

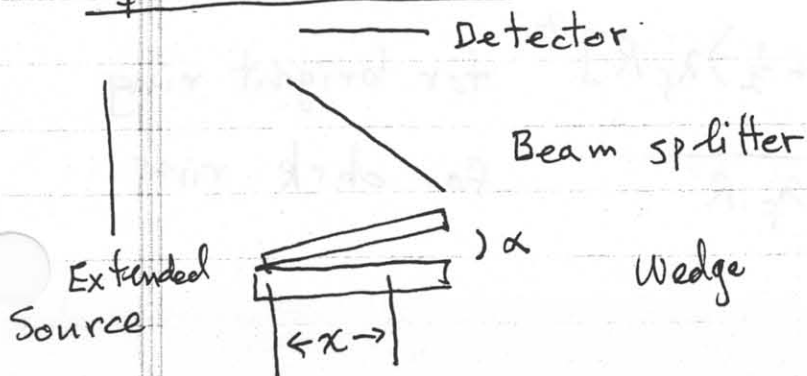
$$\Rightarrow \text{maxima at: } d \cos \theta_t = (2m+1) \left( \frac{\lambda_f}{4} \right)$$

$$m = 0, 1, 2, \dots$$

$$\lambda_f = \frac{\lambda_0}{n_f}$$

Here  $(\theta_i)$  is dominant effect. In other situations,  $(n_f d)$  is dominant.

Equal thickness



View near normal incidence  
 $\Rightarrow$  Fizeau fringes

Fizeau fringes (cont.)

$$\Lambda = 2n_f d \cos(\theta_c)$$

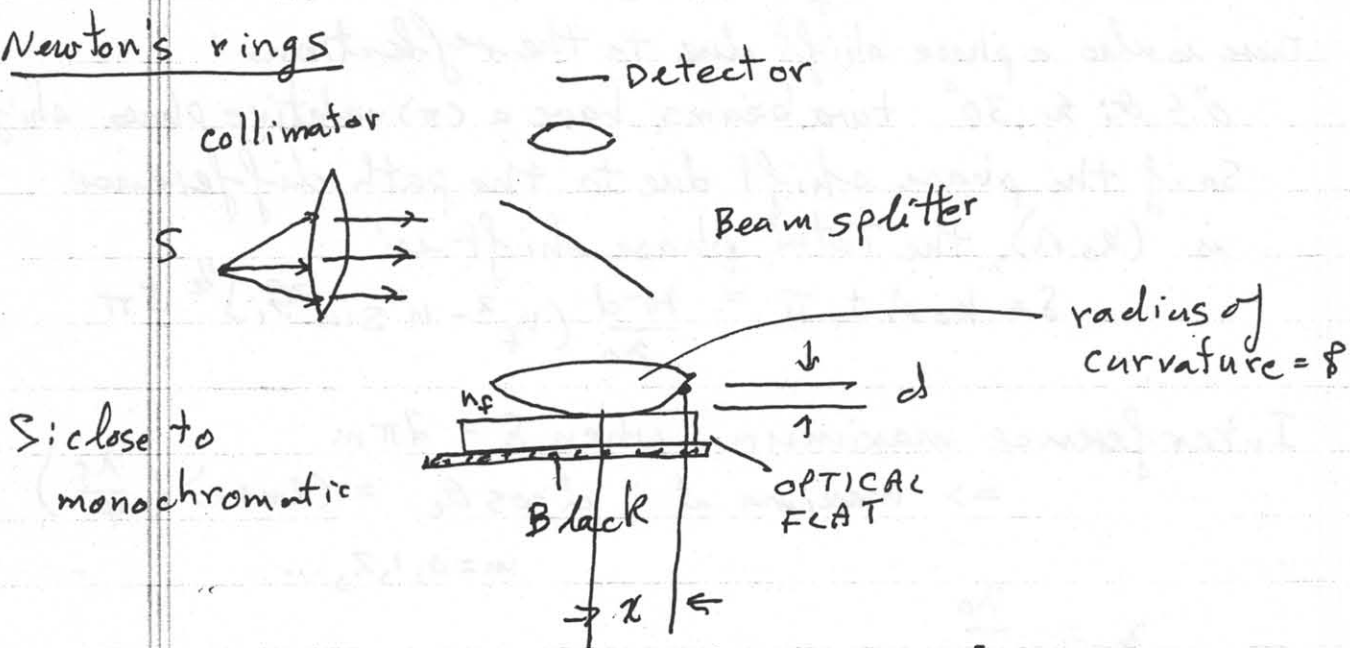
$$d = x\alpha$$

$$\Rightarrow \text{maxima at: } (m + \frac{1}{2})\lambda_0 = 2n_f d_m$$

$$\text{or } (m + \frac{1}{2})\lambda_0 = 2\alpha x_m n_f$$

$$\Rightarrow x_m = \left(\frac{m + \frac{1}{2}}{2\alpha}\right) \lambda_f$$

Newton's rings



S: close to monochromatic

$$x^2 = R^2 - (R-d)^2 = 2Rd - d^2$$

$$R \gg d$$

$$x^2 \approx 2Rd$$

$$\text{Maxima at: } 2n_f d_m = (m + \frac{1}{2})\lambda_0$$

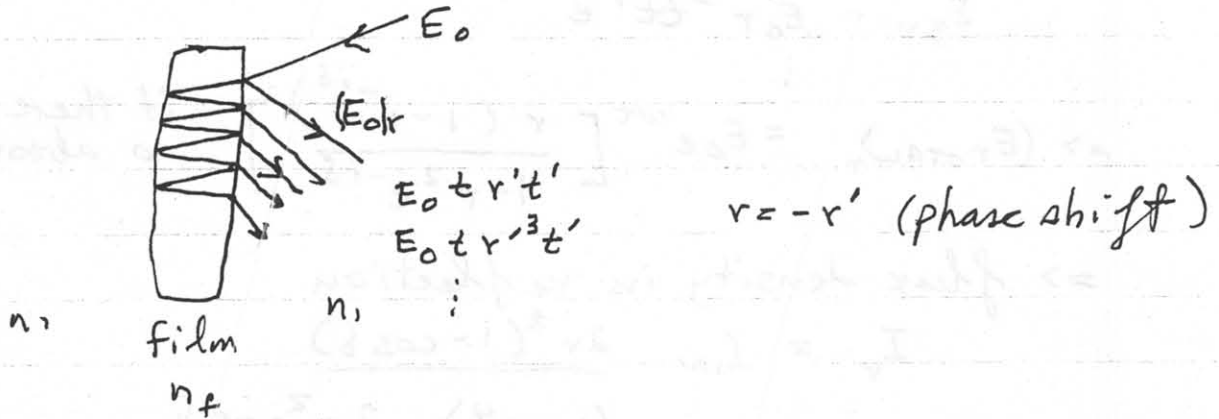
$$\Rightarrow x_m = \left[ (m + \frac{1}{2})\lambda_f R \right]^{1/2} \quad \text{for bright ring}$$

$$= \sqrt{m\lambda_f R} \quad \text{for dark ring}$$

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wedge - sort of - again: Multiple beams -

What if we have more than one beam? Example:



Between adjacent rays:  $\Delta =$  optical path length difference  
 $= 2n_f d \cos \theta_t$

If  $\Delta = m\lambda$ , 2nd, 3rd, 4th, ... rays are in phase, but the 1st ray is  $(\pi)$  out of phase.

$$\begin{aligned} \Rightarrow (E_{TOTAL})_r &= (E_0)_r - (E_0 t r t' + E_0 t r^3 t' + E_0 t r^5 t' + \dots) \\ &= E_0 \left\{ r - t r t' [1 + r^2 + r^4 + \dots] \right\} \\ &= E_0 \left\{ r - \frac{t r t'}{1 - r^2} \right\} \end{aligned}$$

$$t t' = 1 - r^2$$

$$\Rightarrow (E_{TOTAL})_r = 0 \text{ when } \Delta = m\lambda.$$

When  $\Delta = (m + \frac{1}{2})\lambda$ :

1st and 2nd rays are in phase

all other's adjacent rays are  $(\frac{\lambda}{2})$  out of phase

$$\begin{aligned} \Rightarrow (E_{TOTAL})_r &= E_0 r + E_0 r t t' \left\{ 1 - r^2 + r^4 - \dots \right\} \\ &= E_0 r \left[ 1 + \frac{t t'}{1 + r^2} \right] = \frac{2r E_0}{1 + r^2} \end{aligned}$$

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Fizeau + Schlieren

Consider a more general multiple beam interference

$$E_{1r} = E_0 r e^{i\omega t}$$

$$E_{2r} = E_0 r^2 t t' e^{i(\omega t - \delta)}$$

$$E_{3r} = E_0 r^3 t t' e^{i(\omega t - 2\delta)}$$

⋮

$$\Rightarrow (E_{TOTAL})_r = E_0 e^{i\omega t} \left[ \frac{r(1 - e^{-i\delta})}{1 - r^2 e^{-i\delta}} \right] \quad \text{if there is no absorption}$$

$\Rightarrow$  flux density in reflection

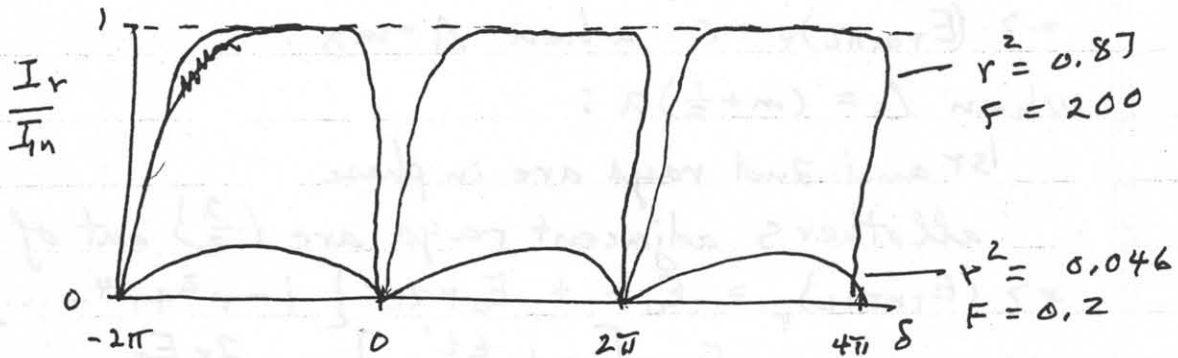
$$I_r = I_{IN} \frac{2r^2(1 - \cos\delta)}{(1 + r^4) - 2r^2 \cos\delta}$$

The minimum for  $I_r$  occurs when  $\delta = 2\pi m$   
and maximum for  $I_r$  occurs when  $\delta = (2m + 1)\pi$

Another term used is the coefficient of finesse

$$F = \left( \frac{2r}{1 - r^2} \right)^2$$

then: 
$$\frac{I_r}{I_{IN}} = \frac{F \sin^2(\delta/2)}{1 + F \sin^2(\delta/2)}$$



Parallel flats



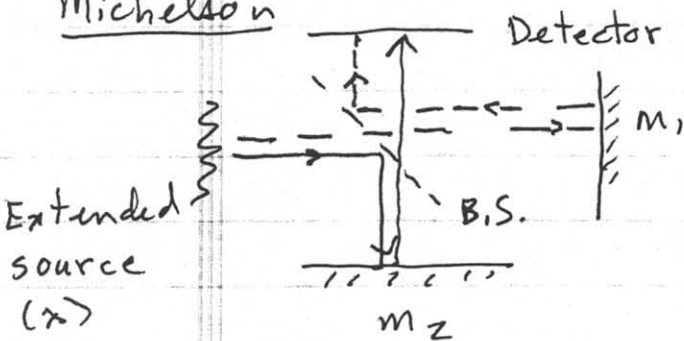
If the flats are parallel to each other and  $s \gg \lambda$  there is a pattern of bull's-eye fringes -- that moves as the eye moves -- as these the fringes are at  $\infty$ .

Source is diffuse.

There are several interferometers that have the same general idea:

- Divide the light into 2 beams, typically by using a  $45^\circ$  beamsplitter.
- Include the perturbation in one beam
- Recombine the two beams, typically by using another  $45^\circ$  beamsplitter
- The perturbation shows up as interference fringes.

Michelson



Main difficulty: Initial alignment to make images of  $m_1$  and  $m_2$  exactly parallel and with exactly zero spacing between them.